

HW #2

Section 1.5 #1, 4a, b)

①

Section 1.6 #6

Section 2.1 #7, 8

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①

$$\frac{d^2 u}{dx^2} + u = 0$$

$$0 < x < L$$

$$u(0) = 0$$

$$u(L) = 0$$

We know the general solution of the

$$\text{ODE is } u(x) = C_1 \cos x + C_2 \sin x$$

$$\text{Since } u(0) = 0 \Rightarrow C_1 = 0$$

$$\text{Thus } u(x) = C \sin x.$$

$$\text{Moreover } 0 = u(L) = C \sin L.$$

$$\text{If } \sin L \neq 0 \Leftrightarrow L \neq n\pi \quad (\text{thus if } L \text{ is not an integer multiple of } \pi)$$

$$n \in \mathbb{Z}$$

then $C = 0$ and $u(x) = 0$ is the only solution.

But if $L = n\pi$ then $u(x) = \sin x$ is another solution and we don't have uniqueness.

④

$$\Delta u = f(x, y, z) \text{ in } D$$

a)

$$\frac{\partial u}{\partial n} = 0 \text{ on } \partial D$$

Let u be a solution. Then if $v = u + C$ where C is an arbitrary constant, then v is another solution since

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Section 1.6 #6

Section 2.1 #7, 8

#1

$$\frac{d^2 u}{dx^2} + u = 0$$

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#4

$$\Delta u = f(x, y, z) \text{ in } D$$

$$\text{a) } \frac{\partial u}{\partial n} = 0 \text{ on } \partial D$$

Let u be a solution. Then if $v = u + C$ where C is an arbitrary constant, then v is another solution since

$$\Delta v = \Delta u = f \quad \text{in } D$$

$$\text{and } \frac{\partial v}{\partial n} = \frac{\partial u}{\partial n} + 0 = 0 \quad \text{on } \partial D$$

No uniqueness

b) Assume u solves $\Delta u = f$ with $\frac{\partial u}{\partial n} = 0$ on ∂D .

$$\text{Then } \iiint_D \Delta u \, dx \, dy \, dz = \iiint_D f \, dx \, dy \, dz$$

$$\begin{aligned} \text{By Gauss Theorem } \iiint_D \Delta u \, dx \, dy \, dz &= \iiint_D \operatorname{div}(\vec{\nabla} u) \, dx \, dy \, dz \\ &= \iint_{\partial D} (\vec{\nabla} u \cdot \vec{n}) \, dS = \iint_{\partial D} \frac{\partial u}{\partial n} \, dS = 0. \end{aligned}$$

Thus a necessary condition for the Neumann problem to have a solution is $\iiint_D f(x, y, z) \, dx \, dy \, dz = 0$.

#6 a) $3u_y + u_{xy} = 0$

We have $a_{11} = a_{22} = 0$ and $a_{12} = \frac{1}{2}$

Since $a_{12}^2 > a_{11}a_{22} = 0$ the PDE is hyperbolic.

b) Let $v = u_y$ then $3v + v_x = 0$ or

$$v_x = -3v. \quad \text{Thus } v(x, y) = f(y) e^{-3x}$$

$$\text{Now } \frac{\partial u}{\partial y}(x, y) = f(y) e^{-3x} \Rightarrow u(x, y) = e^{-3x} \int_0^y f(t) dt + g(x)$$

where f and g are arbitrary functions.

c) Let $u(x, 0) = e^{-3x}$. Then

$$e^{-3x} = g(x) \text{ and thus } u(x, y) = e^{-3x} \left[1 + \int_0^y f(t) dt \right]$$

If in addition $\frac{\partial u}{\partial y}(x, 0) = 0$ we have

$$\frac{\partial u}{\partial y}(x, y) = f(y) e^{-3x} \Rightarrow 0 = \frac{\partial u}{\partial y}(x, 0) = f(0) e^{-3x}$$

$$\Rightarrow f(0) = 0.$$

The solution is not unique. Any function of the form

$$u(x, y) = e^{-3x} \left[1 + \int_0^y f(t) dt \right] \text{ with } f(0) = 0 \text{ is a solution}$$

For example, let $f(y) = y$ then $u(x, y) = e^{-3x} \left(1 + \frac{1}{2} y^2 \right)$ is a solution, but if we take $f(y) = y^2$ then

$$u(x, y) = e^{-3x} \left(1 + \frac{1}{3} y^3 \right) \text{ is also a solution.}$$

#7 The general solution is written

$$u(x, t) = \frac{1}{2} [y(x+ct) + y(x-ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds$$

Assume φ, ψ are odd functions. Then

$$\begin{aligned} u(-x, t) &= \frac{1}{2} [y(-x+ct) + y(-x-ct)] + \frac{1}{2c} \int_{-x-ct}^{-x+ct} \psi(s) ds \\ &= \frac{1}{2} [-y(x-ct) - y(x+ct)] + \frac{1}{2c} \int_{-x-ct}^{-x+ct} \psi(s) ds \end{aligned}$$

If we make the change of variables $s' = -s$ for the last integral we obtain

$$\begin{aligned} u(-x, t) &= -\frac{1}{2} [y(x+ct) + y(x-ct)] - \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(-s') ds' \\ &= -\left\{ \frac{1}{2} [y(x+ct) + y(x-ct)] - \frac{1}{2c} \int_{x+ct}^{x-ct} \psi(s') ds' \right\} \\ &= -\left\{ \frac{1}{2} [y(x+ct) + y(x-ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds \right\} = -u(x, t) \end{aligned}$$

Thus for every t , $u(\cdot, t)$ is also odd.

#8 Consider $u_{tt} = c^2 \left(u_{rr} + \frac{2}{r} u_{rr} \right) = c^2 \Delta u$ (radial Laplacian in 3D)

a) let $v = ru$. Then $v_r = ru_r + u$, $v_{rr} = ru_{rr} + 2u_r$ (*) and $v_{tt} = ru_{tt}$.

Thus $v_{tt} = r u_{tt} = r c^2 \left(u_{rr} + \frac{2}{r} u_r \right) = c^2 \left(r u_{rr} + 2 u_r \right) \stackrel{(*)}{=} \quad (5)$

$$c^2 v_{rr}$$

Thus if u is a solution to $u_{tt} = c^2 \left(u_{rr} + \frac{2}{r} u_r \right)$ and we

set $v = ru$ then v satisfies the 1D Wave Equation.

b) The ^{general} solution to the 1D wave equation is $f(x+ct) + g(x-ct)$ where f and g are arbitrary functions.

Thus $v(r,t) = f(r+ct) + g(r-ct)$ and thus

$$u(r,t) = \frac{1}{r} \left[f(r+ct) + g(r-ct) \right], \text{ Notice that } v(0,t) = 0$$

c) If $u(r,0) = \psi(r)$ and $u_t(r,0) = \phi(r)$ then

$$\psi(r) = \frac{1}{r} (f(r) + g(r)) \rightarrow f(r) + g(r) = r\psi(r)$$

$$\text{and } \phi(r) = u_t(r,0) = \frac{c}{r} [f'(r) - g'(r)]$$

$$\text{Thus } f'(r) + g'(r) = (r\psi(r))' \quad (1)$$

$$f'(r) - g'(r) = \frac{r\phi(r)}{c} \quad (2)$$

$$\text{Thus } f'(r) = \frac{1}{2} \left(\frac{r\psi(r)}{c} + (r\psi(r))' \right)$$

$$g'(r) = \frac{1}{2} \left(\frac{r\psi(r)}{c} - (r\psi(r))' \right)$$

We integrate exactly the same way as when we solved

the 1D wave equation in class and obtain with the

help of $u(r,t) = \frac{1}{r} [f(r+ct) + g(r-ct)]$ that

$$u(r,t) = \frac{1}{2r} \left\{ (r+ct)\varphi(r+ct) + (r-ct)\varphi(r-ct) \right\} + \frac{1}{2cr} \int_{r-ct}^{r+ct} s\psi(s) ds$$

Notice that $v(r,t) = (r+ct)\varphi(r+ct) + (r-ct)\varphi(r-ct) + \frac{1}{2cr} \int_{r-ct}^{r+ct} s\psi(s) ds$

and if φ and ψ are even functions then

$$v(0,t) = ct\varphi(ct) - ct\varphi(-ct) + \frac{1}{2cr} \int_{-ct}^{ct} s\psi(s) ds$$

$$= 0 \quad \text{since } \varphi \text{ is even and } s\psi(s) \text{ is odd.}$$